# Topological Quantum Computation 



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## Topological Quantum Computation



## What is a Quantum Computer?

From [Freedman-Kitaev-Larsen-Wang '03]:
Definition
Quantum Computation is any computational model based upon the theoretical ability to manufacture, manipulate and measure quantum states.


## Prototypical Quantum Computation Scheme

Fix a (classical) function $f: \mathbb{Z}_{2}^{n} \rightarrow \mathbb{Z}_{2}^{n}$.

1. Goal: compute $f(N)$.
2. Encode classical information $N$ as a quantum state $|N\rangle$.
3. Process state: $|N\rangle \rightarrow \sum_{j} a_{j}|j\rangle$ depending on $f$.
4. Measure state: get $|j\rangle$ with probability $\left|a_{j}\right|^{2}$, hopefully $\left|a_{f(N)}\right|^{2} \gg 0$.

## A Universal Quantum Circuit Model

$$
\text { Let } V=\mathbb{C}^{2} \text {. }
$$

Example

- state space ( $n$-qubit): $V^{\otimes n}$
- quantum gate set: $U_{1}:=\left(\begin{array}{llll}1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0\end{array}\right)$,

$$
U_{2}=\left(\begin{array}{cc}
1 & 0 \\
0 & e^{\pi i / 4}
\end{array}\right), U_{3}:=\frac{1}{\sqrt{2}}\left(\begin{array}{cc}
1 & 1 \\
1 & -1
\end{array}\right)
$$

- quantum circuits: $\prod_{j} I_{V}^{\otimes a_{j}} \otimes U_{i j} \otimes I_{V}^{\otimes b_{j}} \in U\left(V^{\otimes n}\right)$, $1 \leq i_{j} \leq 3$.

Theorem
Universal: n-qubit quantum circuits dense in $S U\left(V^{\otimes n}\right)$.

## Remarks on QCM

## Remarks

- Typical physical realization: composite of $n$ identical d-level systems. E.g. $d=2$ : spin- $\frac{1}{2}$ arrays.
- The setting of most quantum algorithms: e.g. Shor's integer factorization algorithm
- Main nemesis: decoherence-errors due to interaction with surrounding material. Requires expensive error-correction...


## Question

How to overcome decoherence? One answer: TOPOLOGY.

## The Braid Group

A key role is played by the braid group $\mathcal{B}_{n}$ generated by $\sigma_{i}$ :


Definition (Artin)
$\mathcal{B}_{n}$ is generated by $\sigma_{i}, i=1, \ldots, n-1$ satisfying:
(R1) $\sigma_{i} \sigma_{i+1} \sigma_{i}=\sigma_{i+1} \sigma_{i} \sigma_{i+1}$
(R2) $\sigma_{i} \sigma_{j}=\sigma_{j} \sigma_{i}$ if $|i-j|>1$

## Anyons

For Point-like particles:

- In $\mathbb{R}^{3}$ : bosons or fermions: $\psi\left(z_{1}, z_{2}\right)= \pm \psi\left(z_{2}, z_{1}\right)$
- Particle exchange $\rightsquigarrow$ reps. of symmetric group $S_{n}$
- In $\mathbb{R}^{2}$ : anyons: $\psi\left(z_{1}, z_{2}\right)=e^{i \theta} \psi\left(z_{2}, z_{1}\right)$
- Particle exchange $\rightsquigarrow$ reps. of braid group $\mathcal{B}_{n}$
- Why? $\pi_{1}\left(\mathbb{R}^{3} \backslash\left\{z_{i}\right\}\right)=1$ but $\pi_{1}\left(\mathbb{R}^{2} \backslash\left\{z_{i}\right\}\right)=F_{n}$ Free group.


$$
C_{1} \not \approx C_{2} \approx C_{3}
$$

## Topological Phases of Matter/Anyons

## Fractional Quantum Hall Liquid



Topological Quantum Computation (TQC) is a computational model built upon systems of topological phases.

## Topological Model



Physics
measure (fusion)
braid anyons
create anyons

## Mathematical Problems

1. Model for anyonic systems/topological phases
2. Classify (models of) topological phases
3. Interpret information-theoretic questions
4. 3-dimensional generalizations?

## Modeling Anyons on Surfaces

Topology of marked surfaces+quantum mechanics"Marks" (boundary components) are labelled by anyons, of which there are finitely many (distinguishable, indecomposable).

## Principle

Superposition: a state is a vector in a Hilbert space $|\psi\rangle \in \mathcal{H}$.
Interpretation


## Principle

The composite state space of two physically separate systems $A$ and $B$ is the tensor product $\mathcal{H}_{A} \otimes \mathcal{H}_{B}$ of their state spaces.

Interpretation


## Principle

Locality: the global state is determined from local information (on disks, plus boundary conditions).

Interpretation
The Hilbert space of a marked surface $M$ is a direct sum over all boundary labelings of a surface $M_{g}$ obtained by cutting $M$ along a circle.


$$
\mathcal{H}=
$$


$\bigoplus_{x} \mathcal{H}_{x}$
( $x^{*}$ is anti-particle to $x$ )

Definition (Nayak, et al '08)
a system is in a topological phase if its low-energy effective field theory is a topological quantum field theory (TQFT).
A 3D TQFT assigns to any surface (+boundary data $\ell$ ) a Hilbert space:

$$
(M, \ell) \rightarrow \mathcal{H}(M, \ell)
$$

Each boundary circle $\bigcirc$ is labelled by $i \in \mathcal{L}$ a finite set of "colors". $0 \in \mathcal{L}$ is neutral. Orientation-reversing map: $x \rightarrow x^{*}$.

## Basic pieces

Any surface can be built from the following basic pieces:

- disk: $\mathcal{H}(\bigcirc ; i)= \begin{cases}\mathbb{C} & i=0 \\ 0 & \text { else }\end{cases}$
- annulus: $\mathcal{H}(\bigcirc ; a, b)= \begin{cases}\mathbb{C} & a=b^{*} \\ 0 & \text { else }\end{cases}$
- pants:



## Two more axioms

Axiom (Disjoint Union)

```
H}[(\mp@subsup{M}{1}{},\mp@subsup{\ell}{1}{})\amalg(\mp@subsup{M}{2}{},\mp@subsup{\ell}{2}{})]=\mathcal{H}(\mp@subsup{M}{1}{},\mp@subsup{\ell}{1}{})\otimes\mathcal{H}(\mp@subsup{M}{2}{},\mp@subsup{\ell}{2}{}
```

Axiom (Gluing)
If $M$ is obtained from gluing two boundary circles of $M_{g}$ together then

$$
\mathcal{H}(M, \ell)=\bigoplus_{x \in \mathcal{L}} \mathcal{H}\left(M_{g}, \ell, x, x^{*}\right)
$$


$(M, \ell)$
$\left(M_{g}, \ell, x, x^{*}\right)$

## Algebraic Part=Modular Category

Morally, $(2+1)$ TQFTs=Modular Categories:
Definition
A modular category $\mathcal{C}$ (over $\mathbb{C}$ ) is
monoidal: $(\otimes, \mathbf{1})$,
semisimple: $X \cong \bigoplus_{i} m_{i} X_{i}$,
linear: $\operatorname{Hom}(X, Y) \in V_{c_{\mathbb{C}}}$,
rigid: $X^{*} \otimes X \mapsto \mathbf{1} \mapsto X \otimes X^{*}$, finite rank: $\operatorname{Irr}(\mathcal{C})=\left\{\mathbf{1}=X_{0}, \ldots, X_{r-1}\right\}$,
spherical: $u_{X} \theta_{X}: X \cong X^{* *}, \operatorname{dim}(X) \in \mathbb{R}$, braided: $c_{X, Y}: X \otimes Y \cong Y \otimes X$,
modular: $\operatorname{Det}\left(\operatorname{Tr}_{\mathcal{C}}\left(c_{X_{i}, X_{j}^{*}} c_{X_{j}^{*}, X_{i}}\right)\right)=\operatorname{Det}\left(S_{i j}\right) \neq 0$.

## Key Data

- fusion rules: $X_{i} \otimes X_{j} \cong \bigoplus_{k} N_{i j}^{k} X_{k}$
- fusion ring representation: $X_{i} \rightarrow N_{i}$ where $\left(N_{i}\right)_{k, j}=N_{i, j}^{k}$
- (modular) S-matrix: $S_{i j}:=\operatorname{Tr}_{\mathcal{C}}\left(c_{X_{i}, X_{j}^{*}} c_{X_{j}^{*}, X_{i}}\right)$
- (Dehn twist) $T$-matrix: $T_{i j}=\delta_{i j} \theta_{i}$
- (quantum)Dimensions: $\operatorname{dim}\left(X_{i}\right):=S_{i 0}$ if unitary, $\operatorname{dim}\left(X_{i}\right)=\max \operatorname{Spec}\left(N_{i}\right)$.


## $(2+1)$ TQFT Anyon Model vs Modular Category

Each axiom has a corresponding physical interpretation:

| TQFT/anyonic system | Category $\mathcal{C}$ |
| :---: | :---: |
| anyon types $x \in \mathcal{L}$ | simple $X$ |
| vacuum $0 \in \mathcal{L}$ | $\mathbf{1}$ |
| $x^{*}$ antiparticle | dual $X^{*}$ |
| $\mathcal{H}(P ; x, y, z)$ state space | Hom $(X \otimes Y, Z)$ |
| particle exchange | braiding $c_{X, X}$ |
| anyon types observable | $\operatorname{det}(S) \neq 0$ |
| topological spin/Dehn twist | $\theta_{X}$ |

## Example: Fibonacci Theory

- $\mathcal{L}=\{0,1\}$
- pants: $\mathcal{H}(P ; a, b, c)= \begin{cases}\mathbb{C} & a=b=c \\ \mathbb{C} & a+b+c \in 2 \mathbb{Z} \\ 0 & \text { else }\end{cases}$
- Define: $V_{k}^{i}:=\mathcal{H}\left(D^{2} \backslash\left\{z_{i}\right\}_{i=1}^{k} ; i, 1, \cdots, 1\right)$
- $\operatorname{dim} V_{n}^{i}= \begin{cases}\operatorname{Fib}(n-2) & i=0 \\ \operatorname{Fib}(n-1) & i=1\end{cases}$


## Classification

## Question (Physics)

How many models exist for a given fixed number of distinguishable indecomposable anyon types?

Theorem (Bruillard,Ng, R, Wang JAMS '15)
There are finitely many modular categories of any given rank $r$.
Proof.
Recall Richard Ng's colloquium at USC on March 12, 2014.
$2 \leq$ rank $\leq 5$ fusion rules (Hong,Ng,Bruillard,Wang,Stong,R.):

| $\|\mathcal{L}\|$ | $\mathcal{C}$ |
| :---: | :---: |
| 2 | $\operatorname{PSU}(2)_{3}, S U(2)_{1}$ |
| 3 | $\mathbb{Z}_{3}, \operatorname{PSU}(2)_{7}, S U(2)_{2}$ |
| 4 | products, $\mathbb{Z}_{4}, \operatorname{PSU}(2)_{9}$ |
| 5 | $\mathbb{Z}_{5}, \operatorname{PSU}(2)_{11}, S U(3)_{4} / \mathbb{Z}_{3}, S U(2)_{4}$ |

## Braid group representations

$\mathcal{B}_{n}$ acts on state spaces:

- Fix anyons $x, y$
- Braid group acts linearly:

$$
\mathcal{B}_{n} \curvearrowright \mathcal{H}\left(D^{2} \backslash\left\{z_{i}\right\} ; x, \cdots, x, y\right)=\operatorname{Hom}\left(X^{\otimes n}, Y\right)
$$



## Universal Anyons

## Question (Quantum Information)

When does an anyon $x$ provide universal computation models?
Basically: when is $\mathcal{B}_{n} \curvearrowright \operatorname{Hom}\left(X^{\otimes n}, Y\right)$ dense?

## Example

Fibonacci $\operatorname{dim}(X)=\frac{1+\sqrt{5}}{2}$ is
universal: braid group $\mathcal{B}_{n}$ image is dense in $S U\left(F_{n}\right) \times S U\left(F_{n-1}\right)$

Example
Ising $\operatorname{dim}(X)=\sqrt{2}$ is not universal: braid group $\mathcal{B}_{n}$ image is a finite group.

## Characterization of Universal anyons

Conjecture ( $\mathrm{R}^{\prime} 07$, property $\mathbf{F}$ )
Anyon $x$ is universal if, and only if, $\operatorname{dim}(X)^{2} \notin \mathbb{Z}$.
Theorem
The property $\mathbf{F}$ conjecture is:

- ( R , Wenzl '14) true for quantum groups
- (Etingof,R,Witherspoon '08) true for group-theoretical categories.


## What do TQCs naturally compute?

Answer
(Approximations to) Link invariants!
Associated to $x \in \mathcal{L}$ is a link invariant $\operatorname{Inv}_{L}(x)$ approximated by the corresponding Topological Model efficiently.


$$
\operatorname{Prob}(\odot) \sim x^{t}\left|\operatorname{lnv}_{l}(\odot)\right|
$$

## Complexity of Jones Polynomial Evaluations

$V_{L}(q)$ Jones polynomial at $q=e^{2 \pi i / \ell}$
Theorem (Vertigan,Freedman-Larsen-Wang)

- (Classical) exact computation of $V_{L}(q)$ at $q=e^{2 \pi i / \ell}$ is: $\left\{\begin{array}{lc}F P & \ell=3,4,6 \\ F P^{\sharp P}-\text { complete } & \text { else }\end{array}\right.$
- (Quantum) approximation of $\left|V_{L}(q)\right|$ at $q=e^{2 \pi i / \ell}$ is $B Q P$.


## 3-dimensional materials

- Point-like particles in $\mathbb{R}^{3}$
- Point-like particles in $\mathbb{R}^{3}$
loop-like particles?



## Two operations: <br> Loop interchange $s_{i}$ : $\bigcirc \leftrightarrow \bigcirc$ and Leapfrogging (read upwards):



$$
\sigma_{i}:
$$




The Loop Braid Group $\mathcal{L B}_{n}$ is generated by
$s_{1}, \ldots, s_{n-1}, \sigma_{1}, \ldots, \sigma_{n-1}$ satisfying:
Braid relations:
(R1) $\sigma_{i} \sigma_{i+1} \sigma_{i}=\sigma_{i+1} \sigma_{i} \sigma_{i+1}$
(R2) $\sigma_{i} \sigma_{j}=\sigma_{j} \sigma_{i}$ if $|i-j|>1$
Symmetric Group relations:
(S1) $s_{i} s_{i+1} s_{i}=s_{i+1} s_{i} s_{i+1}$
(S2) $s_{i} s_{j}=s_{j} s_{i}$ if $|i-j|>1$
(S3) $s_{i}^{2}=1$
Mixed relations:
(M1) $\sigma_{i} \sigma_{i+1} s_{i}=s_{i+1} \sigma_{i} \sigma_{i+1}$
(M2) $s_{i} s_{i+1} \sigma_{i}=\sigma_{i+1} s_{i} s_{i+1}$
(M3) $\sigma_{i} s_{j}=s_{j} \sigma_{i}$ if $|i-j|>1$

## Question

Mathematical models? Do these exist in nature? $(3+1)$ TQFTs...


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## Thank you!

